

On page 4 of Lecture 3(a) we posed the question of determining the range space of the map

$$u \mapsto L(u) = - \int_{t_0}^{t_1} \oint (t_0, \sigma) B(\sigma) u(\sigma) d\sigma.$$

The following answers this.

Lemma: Let $\mathcal{U} = \{u(\cdot): [t_0, t_1] \rightarrow \mathbb{R}^m, u(\cdot) \text{ continuous}\}$

Define the inner product on \mathcal{U}

$$\langle u, v \rangle = \int_{t_0}^{t_1} u(t)^T v(t) dt.$$

Consider the map $L: \mathcal{U} \rightarrow \mathbb{R}^n$

$$L(u) = \int_{t_0}^{t_1} G(t, \sigma) u(\sigma) d\sigma$$

where $G(\cdot)$ is an $n \times m$ continuous matrix-valued function. Giving \mathbb{R}^n the Euclidean inner product, the adjoint map $L^*: \mathbb{R}^n \rightarrow \mathcal{U}$ takes the form

$$\eta \mapsto L^* \eta, \quad (L^* \eta)(t) = G^T(t) \eta$$

and $R(L) = R(LL^*)$.

Proof:

(\Rightarrow) Suppose $x_1 \in R(LL^*)$.

Then $x_1 = LL^*\eta$ where,

$LL^*: \mathbb{R}^n \rightarrow \mathbb{R}^n$ is given by

$$LL^* = \int_{t_0}^{t_1} G_t(\sigma) G_t^\top(\sigma) d\sigma.$$

Let $u_1(\cdot)$ be given by

$$u_1(t) = G_t^\top(t) \eta.$$

$$\text{Then } L(u_1) = \int_{t_0}^{t_1} G_t(\sigma) G_t^\top(\sigma) \eta d\sigma$$

$$= LL^*\eta$$

$$= x_1$$

Hence $x_1 \in R(L)$.

We have shown $R(LL^*) \subseteq R(L)$.

(\Leftarrow) We need to show that if

$x_1 \notin R(LL^*)$ then $x_1 \notin R(L)$.

Given, $x_1 \notin R(LL^*) \exists x_2 \in \text{Ker}(LL^*)^*$

$= \text{Ker}(LL^*)$ (recall LL^* is symmetric), such that $x_1^\top x_2 \neq 0$ [Fredholm alternative].

Suppose $x_1 \in R(L)$ i.e. there exists $u(\cdot)$
such that

$$\int_{t_0}^{t_1} G_1(\sigma) u_1(\sigma) d\sigma = x_1.$$

$$x_1^T x_2 = x_2^T x_1 = \int_{t_0}^{t_1} x_2^T G_1(\sigma) u_1(\sigma) d\sigma$$

$$\neq 0$$

But $0 = x_2^T L^* x_2$

$$= \int_{t_0}^{t_1} x_2^T G_2(\sigma) G_1^T(\sigma) x_2 d\sigma$$

$$= \int_{t_0}^{t_1} \|G_1^T(\sigma) x_2\|^2 d\sigma$$

$$\Rightarrow G_1^T(\sigma) x_2 = 0 \quad \forall \sigma \in [t_0, t_1].$$

$$\Rightarrow x_1^T x_2 = 0, \text{ a contradiction}$$

Hence $x_1 \notin R(L)$. □

Returning to the reachability question, we now conclude,

Theorem

There exists a control $u(\cdot)$ that drives the system

$$\dot{x}(t) = A(t)x(t) + B(t)u(t)$$

$$x(t_0) = x_0$$

to x_1 at time t_1 , iff there is a vector $\eta \in \mathbb{R}^n$ such that

$$W(t_0, t_1)\eta = x_0 - \Phi(t_0, t_1)x_1.$$

In that case, $u(\cdot)$ defined by

$$\begin{aligned} u(t) &= -\left(\Phi(t_0, t)B(t)\right)^T\eta \\ &= -B(t)\Phi(t_0, t)^T\eta \end{aligned}$$

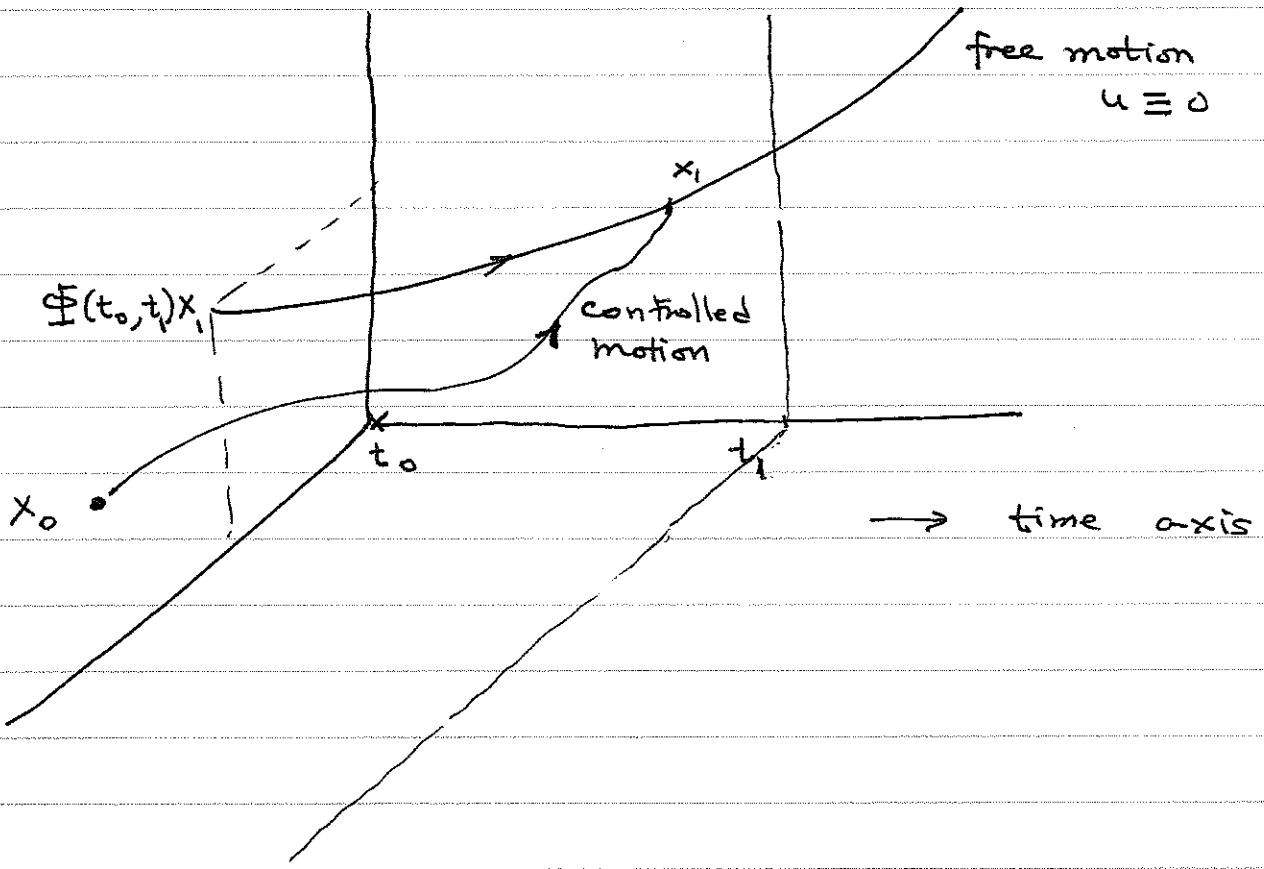
accomplishes the state transfer.

Corollary

If $W(t_0, t_1)$ is invertible then any (x_0, t_0) can be transferred / driven to any (x_1, t_1) , $t_1 > t_0$.

The following diagram (Brockett)

illustrates the situation



The gap between x_0 and $\Phi(t_0, t_1)x_1$, and the uniqueness (given initial conditions) of solutions to the homogeneous equation indicates that free motion cannot transfer (x_0, t_0) to (x_1, t_1) ; only a suitably controlled motion can.

For the setting of time-invariant systems, i.e. $A(t) \equiv A$ a constant and $B(t) \equiv B$ a constant, the reachability Gaussian $W(t_0, t_1)$ can be replaced by an alternative constant matrix with related properties.

Theorem Let $\dot{x} = Ax + Bu$ denote a linear time-invariant system.

Let

$$\begin{aligned} W(t_0, t_1) &= \int_{t_0}^{t_1} e^{(t_0-\sigma)A} BB^T e^{(t_0-\sigma)A^T} d\sigma \\ &= \int_{t_0}^{t_1} e^{-\sigma A} BB^T e^{-\sigma A^T} d\sigma \\ &= W(0, t_1 - t_0) \end{aligned}$$

and let

$$\begin{aligned} W_T &= [B \ A \ B A^2 \ B \dots \ B A^{n-1}] [B \ A \ B A^2 \ B \dots \ B A^{n-1}]^T \\ &= \sum_{k=0}^{n-1} A^k B B^T A^T k \end{aligned}$$

Then, for $t_1 > t_0$

$$R(W(t_0, t_1)) = R(W_T)$$

$$\text{Ker}(W(t_0, t_1)) = \text{Ker}(W_T)$$

(\Rightarrow)

Suppose $x_1 \in \text{Ker}(W(t_0, t_1))$.

Then

$$\begin{aligned} 0 &= x_1^T W(t_0, t_1) x_1, \\ &= \int_{t_0}^{t_1} x_1^T e^{A(t_0-\sigma)} B B^T e^{A^T(t_0-\sigma)} x_1 d\sigma \\ &= \int_{t_0}^{t_1} \|B^T e^{A^T(t_0-\sigma)} x_1\|^2 d\sigma \\ \Rightarrow B^T e^{A^T(t_0-\sigma)} x_1 &= 0 \quad \forall \sigma \in [t_0, t_1]. \end{aligned}$$

Differentiate $(n-1)$ times at $\sigma = t_0$ to obtain,

$$\begin{aligned} B^T x_1 &= 0 \\ B^T A^T x_1 &= 0 \\ \vdots \\ B^T (A^T)^{n-1} x_1 &= 0 \\ \Rightarrow \left(\sum_{k=0}^{n-1} A^k B B^T (A^T)^k \right) x_1 &= 0 \\ \Rightarrow x_1 &\in \text{Ker}(W_T). \end{aligned}$$

(\Leftarrow)

Suppose $x_1 \in \text{Ker}(W_T)$

The Cayley Hamilton theorem says that

$$A^n + p_{n-1} A^{n-1} + p_{n-2} A^{n-2} + \dots + p_0 I = 0$$

where $\chi_A(\lambda) = \lambda^n + p_{n-1} \lambda^{n-1} + p_{n-2} \lambda^{n-2} + \dots + p_0$

is the characteristic polynomial of A .

Using this repeatedly we conclude that

$$e^{tA} = \sum_{i=0}^{n-1} \alpha_i(t) A^i$$

where $\alpha_i(\cdot)$ are functions dependent on the coefficients of the characteristic polynomial.

Then,

$$x_1^T W(t_0, t_1) = \int_{t_0}^{t_1} \sum_{k=0}^{n-1} \alpha_k(t_0 - \tau) x_1^T A^k B^T B A \tau d\tau$$

But $0 = x_1^T w_+$ (by hypothesis)

$$= x_1^T w_+ x_1$$

$$= \sum_{k=0}^{n-1} \| B^T A^k x_1 \|^2$$

$$\Rightarrow x_1^T A^k B = 0 \quad k = 0, 1, 2, \dots, n-1$$

$$\Rightarrow x_1^T W(t_0, t_1) = 0. \quad (\text{substituting in formula above})$$

$$\Rightarrow W(t_0, t_1) x_1 = 0 \quad (\text{since } W(t_0, t_1) = W(t_0, t_1)^T)$$

We have thus shown that

$$\text{Ker}(W(t_0, t_1)) = \text{Ker}(W_T)$$

By Fredholm alternative and the fact that
 $W_T = W_T^T$ and $W(t_0, t_1) = W(t_0, t_1)^T$,
it follows that respective

Kernel and Range
are orthogonal to each other. Hence

$$R(W(t_0, t_1)) = R(W_T).$$

Corollary: $R(W(t_0, t_1)) = R([B \ AB \ \dots \ A^{n-1}B])$

Proof: W_T and $[B \ AB \ \dots \ A^{n-1}B]$ have the same range. \square

Definition We say that a linear constant coefficient system

$$\dot{x} = Ax + Bu$$

is controllable if

$$\text{rank } [B \ AB \ \dots \ A^{n-1}B] = n.$$

Corollary: Controllability \Rightarrow any state can be driven to the origin in finite time.

Theorem

Consider a single-input time-invariant linear system of the form

$$\dot{x} = Ax + bu$$

where $b \in \mathbb{R}^n$ A is $n \times n$. Then there is choice of basis in state space such that in this new basis the system can be brought to the canonical form

$$\dot{z} = A_c z + b_c u$$

where

$$A_c = \begin{bmatrix} 0 & 1 & 0 & \cdots & 0 \\ 0 & 0 & 1 & \cdots & 0 \\ 0 & 0 & 0 & \cdots & 1 \\ -p_0 & -p_1 & \cdots & -p_{n-1} & \end{bmatrix}$$

$$b_c = \begin{bmatrix} 0 \\ 0 \\ \vdots \\ 1 \end{bmatrix}$$

where $\chi_A(s) = s^n + p_{n-1}s^{n-1} + p_{n-2}s^{n-2} + \cdots + p_1s + p_0$.

Proof

First note that a change of basis yields
change of coordinates

$$\mathbf{z} = P\mathbf{x},$$

where P is a nonsingular matrix and

$$\dot{\mathbf{z}} = P\dot{\mathbf{x}}$$

$$= P(\mathbf{A}\mathbf{x} + \mathbf{b}\mathbf{u})$$

$$= PAP^{-1}\mathbf{z} + Pb\mathbf{u}$$

$$= A_c\mathbf{z} + b_c\mathbf{u}$$

Thus we seek P such that

$$A_c = PAP^{-1}$$

$$b_c = Pb.$$

By controllability

$$\text{rank } [b, Ab, \dots, A^{n-1}b] = n$$

i.e. the vectors $\{b, Ab, \dots, A^{n-1}b\}$

constitute a linearly independent set. Now consider the 'triangular' linear combinations of the form,

$$v_1 = A^{n-1}b + p_{n-1}A^{n-2}b + \dots + p_1b$$

$$v_2 = A^{n-2}b + p_{n-1}A^{n-3}b + \dots + p_2b$$

$$v_3 = A^{n-3}b + p_{n-1}A^{n-4}b + \dots + p_3b$$

$$v_{n-1} = Ab + p_{n-1}b$$

$$v_n = b$$

The set of vectors $\{v_1, v_2, \dots, v_n\}$ is clearly linearly independent. (OBSERVE: Proceeding from v_n , a new direction $A^k b$ is added each time we move up the ladder through v_{n-1}, v_{n-2}, \dots , to v_1 .)

Let us express the matrix A and the matrix b in this new basis.

Clearly, since u takes values in \mathbb{R}^l choosing ' 1 ' to be basis for \mathbb{R}^l , we see

$$b \cdot 1 = \sum_{i=1}^n \tilde{b}_i v_i$$

$$= 0 \cdot v_1 + 0 \cdot v_2 + \dots + 0 \cdot v_{n-1} + 1 \cdot v_n$$

Thus b has matrix representation

$$b_c = \begin{pmatrix} 0 \\ 0 \\ \vdots \\ 1 \end{pmatrix}$$

in the new basis.

By Cayley Hamilton theorem,

$$A^n = -p_{n-1} A^{n-1} - p_{n-2} A^{n-2} \dots - p_0 I$$

$$\begin{aligned} A v_1 &= A^n b + p_{n-1} A^{n-1} b + \dots + p_1 A b \cancel{+ p_0 b} \\ &= -p_0 b \\ &= -p_0 v_n \end{aligned}$$

$$\begin{aligned} A v_2 &= A^{n-1} b + p_{n-1} A^{n-2} b + p_{n-2} A^{n-3} b + \dots + p_2 A b \\ &= v_1 - p_1 b \\ &= v_1 - p_1 v_n \end{aligned}$$

$$\begin{aligned} A v_3 &= A^{n-2} b + p_{n-1} A^{n-3} b + \dots + p_3 A b \\ &= v_2 - p_2 b \\ &= v_2 - p_2 v_n \end{aligned}$$

$$\begin{aligned} A v_{n-1} &= A^2 b + p_{n-1} A b \\ &= v_{n-2} - p_{n-2} b \\ &= v_{n-2} - p_{n-2} v_n \end{aligned}$$

$$\begin{aligned} A v_n &= A b \\ &= v_{n-1} - p_{n-1} b \\ &= v_{n-1} - p_{n-1} v_n \end{aligned}$$

It follows that the matrix representation of the linear map $A : \mathbb{R}^n \rightarrow \mathbb{R}^n$ in the new basis $\{v_1, v_2, \dots, v_n\}$ is given by,

$$A_c = \begin{bmatrix} 0 & 1 & 0 & \cdots & 0 \\ 0 & 0 & 1 & \cdots & 0 \\ \vdots & & & & \vdots \\ 0 & 0 & 0 & \cdots & 1 \\ -p_0 & -p_1 & -p_2 & \cdots & -p_{n-1} \end{bmatrix}.$$

This completes the proof.

(We did not write down P explicitly. Instead we invented a basis and expressed A as A_c in this new basis, and similarly b as b_c in this new basis). With $\{e_1, e_2, \dots, e_n\}$ as the standard basis in \mathbb{R}^n , the new basis $\{v_1, v_2, \dots, v_n\}$ can be expressed as the matrix

$$[v_1, v_2, \dots, v_n] = [b \ A b \ A^2 b \ \cdots \ A^{n-1} b] \begin{bmatrix} p_0 & p_1 & p_2 & p_3 & \cdots & p_{n-1} \\ p_1 & p_2 & p_3 & p_4 & \cdots & \vdots \\ p_2 & p_3 & p_4 & p_5 & \cdots & \vdots \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ p_{n-1} & p_0 & p_1 & p_2 & \cdots & p_{n-2} \\ 1 & & & & & \end{bmatrix}$$

$$= T$$

the change of basis from $\{e_1, \dots, e_n\}$ to $\{v_1, v_2, \dots, v_n\}$

Adopting the convention introduced in
Lecture 1 (a), page 10,

$$P = T^{-1}$$

Verify these steps.)

Remark

The pair $[A_c, b_c]$ or equivalently
the system

$$\dot{z} = A_c z + b_c u$$

constitute the control canonical
form. Observe that, if we
let

$$y(t) = z_1(t),$$

then

$$\frac{d}{dt^n} y(t) + p_{n-1} \frac{d}{dt^{n-1}} y(t) + \dots + p_1 \frac{dy}{dt} + p_0 y$$
$$= u(t)$$

Thus the control canonical form reveals
structure of an n^{th} order differential
equation underlying the original system.